Thermodynamic binding networks

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DNA XARAAAAA







4/41



4/41

short





Bind



Bind



Bind





Displace

















Leak in strand displacement experiments



Source: Lulu Qian, Erik Winfree. Scaling Up Digital Circuit Computation Science 332, 2011

Leak in strand displacement experiments



Reducing Leak



[Boya Wang, Chris Thachuk, Andrew Ellington, David Soloveichik. *The Design Space of Strand Displacement Cascades with Toehold-Size Clamps* <u>DNA Computing Conference</u>, 2017]



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What causes leak "thermodynamically"?


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less favorable

more favorable













Geometry-Free Model:

The domains within a monomer are unordered

Monomer = collection of domains Configuration = how monomers are bound

Energetic favorability: Bonds and complexes

all else equal, more bonds = more favorable







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energy = w*#bonds + #complexes

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- We often consider a natural limiting case:
 - favoring # bonds <u>infinitely</u> over #complexes
 - require maximal #bonds formed; use #complexes only as tiebreaker
 - Corresponds to bonds that are so strong they <u>cannot spontaneously dissociate</u>, but can <u>exchange</u> with each other to find configurations with more complexes



saturated = maximum #bonds formed

Saturated

b

b







If we're careful to make starred binding sites limiting, then saturated = all starred sites are bound

Computing via Thermodynamic Equilibrium

















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- Output convention?
 - Obvious one: "there's a unique stable configuration with the correct output"
 - It's problematic, so we have a one-sided convention:
 - if correct output is 0, unique stable configuration with correct answer
 - if correct output is 1, then both the "output=1" and "output=0" configurations are stable

Composable AND gate with entropy gap 3



Rather than release a single output monomer, it suffices to gather all output domains on one complex.
Kinetic pathways and energy barriers

Pathways

Thermodynamics: Which configurations are energetically favorable **Kinetics**: How a system moves between configurations over time





$X \longrightarrow Y + Z$









. . .

h_z*

h_v*

$X \longrightarrow Y + Z$







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$\phi \rightarrow Y + Z$



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What causes leak "kinetically"?

 $\phi \longrightarrow \mathbf{Y} + \mathbf{Z}$









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Since all saturated configurations have an equal number of bonds, we can focus solely on the number of complexes

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$x_{11} x_{12} x_{13} x_{14}$	x ₁₁ *	x ₁₂ *	x ₁₃ *	x ₁₄ *	x ₁₁	x ₁₂	х ₁₃	x ₁₄
x ₂₁ x ₂₂ x ₂₃ x ₂₄	x ₂₁ *	x ₂₂ *	x ₂₃ *	x ₂₄ *	x ₂₁	x ₂₂	x ₂₃	x ₂₄
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x ₄₁ x ₄₂ x ₄₃ x ₄₄	x ₄₁ *	x ₄₂ *	x ₄₃ *	x ₄₄ *	x ₄₁	x ₄₂	х ₄₃	x ₄₄



















progress





30




































X₁₁

X21



X₄₃

^44

X₄₂

progress

(X₄₁)

(X₄₂)



31

x₂₁

X₃₁

. X₄₁



X₁₂

X₂₂

X₃₂

(X₄₂)

x₂₁

X₃₁

x₄₁

































































x ₁₂	x ₁₃	x ₁₄
x ₂₂	x ₂₃	x ₂₄
х ₃₂	х ₃₃	х ₃₄
x ₄₂	х ₄₃	х ₄₄
	x ₁₂ x ₂₂ x ₃₂ x ₄₂	x ₁₂ x ₁₃ x ₂₂ x ₂₃ x ₃₂ x ₃₃ x ₄₂ x ₄₃





























Along a catalyzed pathway, the barrier is 1



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- Unknown for next prime power *n*=10:
 - trivial upper bound is 11 days
 - best known lower bound is 3




















Thermodynamic self-assembly

Grafting the TBN model onto self-assembly

A modest goal

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 - <u>size of a complex</u> = # monomers in the complex
- Formal: Design a set of monomer types so that, for all *S* ∈ N, there is a stable complex of size at least *S*.
- Easy to do in Abstract Tile Assembly Model:









attempt 2:





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D,M = O(1), S = arbitrarily large



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 $S \approx D^2$

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$$c_1 d_1^* - d_1 c_1 d_2^* - d_2 c_1 d_3^* - d_3 c_1$$

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$$S \approx 2^D$$
?







Stable complexes have at most exponential size

Theorem: Any thermodynamic binding network with

- D domain types,
- *M* monomer types,
- $\leq A$ domains per monomer type

(note $D/A \leq M \leq A^{D+1}$)

Has stable complexes of size $\leq 2(M+D)(AD)^{2D+3} = \text{poly}(D^D)$ if A = O(1)



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- If some path has > 2D edges, it must repeat some ordered pair (d_i,d_i*) or (d_i*,d_i)
- Break into two saturated complexes as shown.



Monomers as vectors

• monomer {a, b*,b*, d,d,d,d*, e,e*} represented as (1,-2,0,3,0)

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- sum of many monomers gives the number of excess domains in a fully bound (saturated) complex with those monomers
 - i.e., 2 copies of above monomer 2·(1,-2,0,3,0) = (2,-4,0,6,0) have an excess of 2 a's, 4 b*'s, 0 c's, 6 d's, 0 e's
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- 6. i.e., we can split S_2 into 2 disjoint nonempty nonnegative subsets, S_1 and $M \cdot d$. **QED**

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If Ax = b has a solution, it has a "small" solution... $\max_i x_i \le \exp(\max_{ij}(A_{ij}, b_j))$

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<u>Question</u>: is there a nonnegative integer vector **x** such that **Ax** = **b**?

- 0/1-INTEGER-PROGRAMMING is NP-complete (Karp 1972).
- <u>Non-obvious fact</u>: INTEGER-PROGRAMMING is in NP. (independently due to [Borosh and Treybig 1976], [Gathen and Sieveking 1978], [Kannan and Monma 1978])
 If Ax = b bas a solution, it bas a "small" solution, max x < oxp(max (A, b))

If **Ax** = **b** has a solution, it has a "small" solution... max_i **x**_i ≤ exp(max_{ij}(**A**_{ij},**b**_j))

- Papadimitriou's proof: [On the complexity of integer programming. Papadimitriou, JACM 1981]
 - If **x** is a *large enough* solution, there is 0 < y < x, $y \in \mathbb{N}^m$, such that Ay = 0.

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 - Defining $\mathbf{z} = \mathbf{x} \mathbf{y}$, $\mathbf{A}\mathbf{z} = \mathbf{A}(\mathbf{x} \mathbf{y}) = \mathbf{A}\mathbf{x} \mathbf{A}\mathbf{y} = \mathbf{A}\mathbf{x} \mathbf{0} = \mathbf{b}$.

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 - Defining z = x y, Az = A(x y) = Ax Ay = Ax 0 = b.
 - So z is a strictly smaller solution than x: x cannot be the *smallest* solution.

Farkas' Lemma

Given vectors \mathbf{m}_1 , \mathbf{m}_2 , ..., they obey one of two constraints:

a) are directions of balanced forces b) lie on one side of some hyperplane



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- So whatever bonds were broken to separate y can be <u>re-bound within y</u>.
- By symmetry, the same bonds in z = c y can be rebound within z.

monomer collection $\mathbf{c} \in \mathbb{N}^{M}$



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• Then "small" $\geq -\sum_{i \in S} (\mathbf{h} \cdot \mathbf{m}_i) \mathbf{c}(i) = \sum_{i \notin S} (\mathbf{h} \cdot \mathbf{m}_i) \mathbf{c}(i) \geq \sum_{i \notin S} \mathbf{c}(i)$ **c**(i) (count of i'th monomer) is above since $\mathbf{h} \cdot \mathbf{m}_i \geq 1$ small by definition, and $\mathbf{h} \cdot \mathbf{m}_i = O(1)$

Applying thermodynamic model to tile assembly

- Let's incorporate the thermodynamic binding network model into the abstract tile assembly model.
- How can we create a large assembly from a small number of tile types?








С

С

С

s₀₁

C

0*



С

С

С

n

0

С

С

Difference is that each row (corresponding to bits of the same significance) has glues labeled with the row number



Conclusions

- Strong bonds (surprisingly) aren't sufficient to self-assemble large thermodynamically stable structures. *Geometry helps*!
- Kinetically self-assembling a thermodynamically stable structure has very strong guarantees on errors:
 - target structure eventually results despite arbitrary kinetic errors.
 - If it's the only stable structure, and free energy of other structures is much less, then it's the only result you'll see.
- Bad news: **NP**-complete to tell if a given configuration is unstable... even **NP**-hard to approximate entropy of stable configuration:

[Breik, Thachuk, Heule, Soloveichik, *Computing properties of stable configurations of thermodynamic binding networks*, <u>Theoretical Computer Science</u> 2019]