# On the shift-invert Lanczos method for the buckling eigenvalue problem 

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#### Abstract

We consider the problem of extracting a few desired eigenpairs of the buckling eigenvalue problem $K x=\lambda K_{G} x$, where $K$ is symmetric positive semi-definite, $K_{G}$ is symmetric indefinite, and the pencil $K-\lambda K_{G}$ is singular, namely, $K$ and $K_{G}$ share a nontrivial common nullspace. Moreover, in practical buckling analysis of structures, bases for the nullspace of $K$ and the common nullspace of $K$ and $K_{G}$ are available. There are two open issues for developing an industrial strength shift-invert Lanczos method: (1) the shift-invert operator $\left(K-\sigma K_{G}\right)^{-1}$ does not exist or is extremely ill-conditioned, and (2) the use of the semi-inner product induced by $K$ drives the Lanczos vectors rapidly toward the nullspace of $K$, which leads to a rapid growth of the Lanczos vectors in norms and causes permanent loss of information and the failure of the method. In this paper, we address these two issues by proposing a generalized buckling spectral transformation of the singular pencil $K-\lambda K_{G}$ and a regularization of the inner product via a low-rank updating of the semi-positive definiteness of $K$. The efficacy of our approach is demonstrated by numerical examples, including one from industrial buckling analysis.


## KEYWORDS

buckling analysis, eigenvalue problem, Lanczos method, shift-invert, singular pencil

## 1 | INTRODUCTION

We consider the buckling eigenvalue problem

$$
\begin{equation*}
K x=\lambda K_{G} x, \tag{1}
\end{equation*}
$$

where $K$ and $K_{G}$ are $n \times n$ sparse symmetric matrices, and $K$ is positive semi-definite and $K_{G}$ is indefinite. Furthermore, the pencil $K-\lambda K_{G}$ is singular, that is, the matrices $K$ and $K_{G}$ share a nontrivial common nullspace $\mathcal{Z}_{c}$. We are interested in (i) extracting a few nonzero finite eigenvalues around a prescribed shift $\sigma \neq 0$ and the associated eigenvectors $x$ perpendicular to the common nullspace $\mathcal{Z}_{c}$, and (ii) counting the number of eigenvalues of $K-\lambda K_{G}$ in a given interval ( $\alpha, \beta$ ). As in practical buckling analysis of structures, we assume that a basis $Z \equiv\left[Z_{N} Z_{C}\right]$ of the nullspace of $K$ and a basis $Z_{C}$ of the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$ are available, and the pencil $K-\lambda K_{G}$ is simultaneously diagonalizable.

The buckling eigenvalue problem (1) arises from the buckling analysis in structural engineering, where $K$ is referred to as the stiffness matrix and $K_{G}$ is referred to as the geometric stiffness matrix. The eigenvalue $\lambda$ is used to determine
the critical load at which a structure may become unstable (Reference 1, p. 433), and the eigenvector $x$ is the associated buckling shape. The bases for the nullspace of $K$ and the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$ can be extracted from the algebraic or geometric structure of the problem. ${ }^{1-3}$

The buckling eigenvalue problem (1) remains an outstanding computational challenge in numerical linear algebra ${ }^{4,5}$ and in industrial applications. ${ }^{6}$ When the pencil $K-\lambda K_{G}$ is regular and $K$ is positive definite, a common practice for computing eigenpairs around a given shift $\sigma$ is to convert (1) into the following ordinary eigenproblem via a so-called buckling spectral transformation

$$
\begin{equation*}
\left(K-\sigma K_{G}\right)^{-1} K x=\frac{\lambda}{\lambda-\sigma} x, \tag{2}
\end{equation*}
$$

see References ${ }^{7-10}$. Since $\left(K-\sigma K_{G}\right)^{-1} K$ is symmetric with respect to $K$, the Lanczos method with $K$-inner product can be immediately used to solve the eigenproblem (2). This approach is referred to as the shift-invert Lanczos method and has been widely used, including in a number of industrial strength eigensolvers, such as LS-DYNA. ${ }^{6}$

However, when $K$ is positive semi-definite and $K-\lambda K_{G}$ is singular, we have the following two issues:

1. Since the pencil $K-\lambda K_{G}$ is singular or near singular, that is, the matrices $K$ and $K_{G}$ share a nontrivial common nullspace $\mathcal{Z}_{c}$, the shift-invert matrix $\left(K-\sigma K_{G}\right)^{-1}$ does not exist or is extremely ill-conditioned.
2. Since the matrix $K$ is positive semi-definite, the inner product induced by $K$ causes the Lanczos vectors driven rapidly toward the nullspace of $K .{ }^{4,5,8,11}$ It results in the large norms of the Lanczos vectors, which introduces large rounding errors. The accuracy of the computed solutions degrades and the procedure can even fail.

These issues have been studied since the early development of the shift-invert Lanczos method in the 1980s. Nour-Omid et al. ${ }^{8}$ proposed a modified formulation of the Ritz vectors to refine the computed solutions. Meerbergen ${ }^{4}$ proposed to control the norms of the Lanczos vectors by applying implicit restart. ${ }^{12}$ More recently, Stewart ${ }^{5}$ gave a detailed analysis to show that the loss of information caused by the growth of the Lanczos vectors is permanent.

One way to address the issues is to consider constraints on the degrees of freedom (Reference 1, p. 272). By removing the redundant degrees of freedom, the buckling eigenvalue problem (1) can be transformed into an equivalent symmetric definite generalized eigenvalue problem. This approach, however, could significantly increase the number of nonzero entries in the shifted matrix, leading to huge amount of memory for the factorization. The constraints can also be imposed by augmenting (1) using the Lagrange formulation. ${ }^{6}$ But both the augmented matrices become indefinite and the shift-invert Lanczos method is not applicable. Alternative way is to enforce the Lanczos vectors in the desired subspace by deflation. ${ }^{8}$ Still, the stability could be a concern.

In this paper, we address the two issues by first proposing a generalized buckling spectral transformation of the singular pencil $K-\lambda K_{G}$, and a regularization of the inner product via a low-rank updating of the positive semi-definite matrix $K$. Then a shift-invert Lanczos method for the buckling eigenvalue problem (1) is developed. We will discuss implementation of the matrix-vector product for the computational kernel of the shift-invert Lanczos method, and propose way to count the number of eigenvalues in a given interval $(\alpha, \beta)$ for validation.

The rest of the paper is organized as follows. In Section 2, we first present a canonical form of the pencil $K-\lambda K_{G}$, and propose a generalized buckling spectral transformation, and a regularization of the inner product. In Section 3, we discuss the implementation of the shift-invert Lanczos method with the generalized buckling spectral transformation and the regularized inner product. In Section 4, we discuss way to count the number of eigenvalues in an interval. Efficacy of the proposed approach is demonstrated in Section 5. Concluding remarks are given in Section 6.

Following the convention of matrix computations, we use the upper case letters for matrices and lower case letters for vectors. In particular, we use $I_{n}$ for the identity matrix of dimension $n$ with $e_{j}$ being the $j$ th column. If not specified, the dimensions of matrices and vectors conform to the dimensions used in the context. ${ }^{T}$ is for transpose,.$^{\dagger}$ for pseudo-inverse, $\|\cdot\|_{1}$ for 1-norm, and $\|\cdot\|_{2}$ and $\|\cdot\|_{F}$ for 2-norm and Frobenius norm, respectively. Also, $\kappa_{2}(\cdot)$ is for the 2-norm condition number. We use $A^{-T}$ for the inverse of the matrix $A^{T}$. The range and the nullspace of a matrix $A$ are denoted by $\mathcal{R}(A)$ and $\mathcal{N}(A)$, respectively. The direct sum of two subspaces $S_{1}$ and $S_{2}$ is denoted by $S_{1} \oplus S_{2}$. The orthogonal complement to a subspace $S$ is denoted by $S^{\perp}$ and the orthogonal projection onto a subspace $S$ is denoted by $\mathcal{P}_{S}$. $v_{+}(S), v_{-}(S)$ and $v_{0}(S)$ denote the numbers of positive, negative and zero eigenvalues of a symmetric matrix $S$, respectively. Other notations will be explained as used.

## 2 | THEORY

## 2.1 | Canonical form

We start with a canonical form of the pencil $K-\lambda K_{G}$. For the compactness of presentation, we interchange the roles of $K$ and $K_{G}$ in (1) and consider the reversal of the pencil $K-\lambda K_{G}$, that is, $K_{G}-\lambda^{\#} K$.

Theorem 1. For the pencil $K_{G}-\lambda^{\#} K$, there exists a nonsingular matrix $W \in \mathbb{R}^{n \times n}$ such that

$$
W^{T} K_{G} W=\begin{gather*}
n_{2}  \tag{3}\\
n_{1}
\end{gather*}\left[\begin{array}{lll}
n_{1} & n_{2} & n_{3} \\
\Lambda_{1}^{\#} & & \\
& \Lambda_{2}^{\#} & \\
& & 0
\end{array}\right] \text { and } W^{T} K W=\begin{array}{rcc}
n_{1} & n_{2} & n_{3} \\
n_{2} \\
n_{3}
\end{array}\left[\begin{array}{lll}
I_{n_{1}} & & \\
& 0 & \\
& & 0
\end{array}\right],
$$

where $\Lambda_{1}^{\#}$ and $\Lambda_{2}^{\#}$ are diagonal matrices with real diagonal entries, and $\Lambda_{2}^{\#}$ is nonsingular. Furthermore, by conformally partitioning $W=\left[W_{1}, W_{2}, W_{3}\right]$, we have

$$
\begin{equation*}
W_{3}^{T} W_{1}=0 \quad \text { and } \quad W_{3}^{T} W_{2}=0 \tag{4}
\end{equation*}
$$

Proof. see Appendix A.
Remark 1. By the canonical form (3), we immediately know that (i) the columns of $W_{3}$ span the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$, and the columns of [ $W_{1} W_{2}$ ] span the orthogonal complement to $\mathcal{Z}_{c}$, that is, $\mathcal{Z}_{c}^{\perp}$; (ii) the columns of $W_{1}$ are eigenvectors associated with real finite eigenvalues $\left(\Lambda_{1}^{\#}, I_{n_{1}}\right)$ of the pencil $K_{G}-\lambda^{\#} K$ and are perpendicular to $\mathcal{Z}_{c}$; (iii) The columns of $W_{2}$ are eigenvectors associated with an infinite eigenvalue ( $\Lambda_{2}^{\#}, 0$ ) of the pencil $K_{G}-\lambda^{\#} K$ and are perpendicular to $\mathcal{Z}_{c}$; (iv) For $x \in \mathcal{Z}_{c},\left(\lambda^{\#}, x\right)$ is an eigenpair of the pencil $K_{G}-\lambda^{\#} K$ for any $\lambda^{\#} \in \mathbb{C}$.

## 2.2 | Generalized buckling spectral transformation

Mathematically, a generalized buckling spectral transformation of the singular pencil $K-\lambda K_{G}$ is to replace the inverse in (2) by the pseudo-inverse and leads to the ordinary eigenvalue problem

$$
\begin{equation*}
C x=\mu x \quad \text { with } \quad C=\left(K-\sigma K_{G}\right)^{\dagger} K \tag{5}
\end{equation*}
$$

where $\left(K-\sigma K_{G}\right)^{\dagger}$ is the pseudo-inverse of the singular matrix $K-\sigma K_{G}$ (Reference 13, p.290). Note that the nonzero real shift $\sigma$ cannot be an eigenvalue of the pencil $K-\lambda K_{G}$.

We now present the relationship of nontrivial eigenpairs between the original buckling eigenvalue problem (1) and the ordinary eigenvalue problem (5). We first use the canonical form (3) to derive an eigenvalue decomposition of $C$ and provide the eigenvalue and eigenvector relations between $C$ and $K_{G}-\lambda^{\#} K$.

Lemma 1. With the canonical form (3) in Theorem 1, an eigenvalue decomposition of the matrix C defined in (5) is given by

$$
C W=W\left[\begin{array}{lll}
\left(I_{n_{1}}-\sigma \Lambda_{1}^{\#}\right)^{-1} & &  \tag{6}\\
& 0 & \\
& & 0
\end{array}\right] .
$$

Proof. Recall that, since the matrix $K-\sigma K_{G}$ is symmetric,

$$
\begin{equation*}
\mathcal{R}\left(K-\sigma K_{G}\right)=\mathcal{N}\left(K-\sigma K_{G}\right)^{\perp}=\mathcal{Z}_{c}^{\perp} \tag{7}
\end{equation*}
$$

In addition, by the condition (4) in the canonical form (3), we have

$$
\begin{equation*}
\mathcal{R}\left(W_{1}\right) \oplus \mathcal{R}\left(W_{2}\right)=\mathcal{R}\left(W_{3}\right)^{\perp}=\mathcal{Z}_{c}^{\perp} . \tag{8}
\end{equation*}
$$

Therefore, from (7) and (8),

$$
\begin{equation*}
\mathcal{R}\left(K-\sigma K_{G}\right)=\mathcal{R}\left(W_{1}\right) \oplus \mathcal{R}\left(W_{2}\right)=\mathcal{R}\left(W_{3}\right)^{\perp}=\mathcal{Z}_{c}^{\perp} . \tag{9}
\end{equation*}
$$

Now note that, from the canonical form (3),

$$
W^{T} K W=\left[\begin{array}{lll}
I_{n_{1}} & &  \tag{10}\\
& 0 & \\
& & 0
\end{array}\right] \text { and } \quad W^{T}\left(K-\sigma K_{G}\right) W=\left[\begin{array}{lll}
I_{n_{1}}-\sigma \Lambda_{1}^{\#} & & \\
& -\sigma \Lambda_{2}^{\#} & \\
& & 0
\end{array}\right] .
$$

Therefore, we have

$$
W^{T} K W=\left[\begin{array}{lll}
I_{n_{1}} & &  \tag{11}\\
& 0 & \\
& & 0
\end{array}\right]=W^{T}\left(K-\sigma K_{G}\right) W\left[\begin{array}{lll}
\left(I_{n_{1}}-\sigma \Lambda_{1}^{\#}\right)^{-1} & & \\
& 0 & \\
& & 0
\end{array}\right] .
$$

Left multiplying (11) by $\left(K-\sigma K_{G}\right)^{\dagger} W^{-T}$,

$$
\left(K-\sigma K_{G}\right)^{\dagger} K W=\left(K-\sigma K_{G}\right)^{\dagger}\left(K-\sigma K_{G}\right) W\left[\begin{array}{lll}
\left(I_{n_{1}}-\sigma \Lambda_{1}^{\#}\right)^{-1} & &  \tag{12}\\
& 0 & \\
& & 0
\end{array}\right] .
$$

The pseudo-inverse $\left(K-\sigma K_{G}\right)^{\dagger}$ satisfies the Moore-Penrose conditions (Reference 13, p. 290), which give

$$
\begin{equation*}
\left(K-\sigma K_{G}\right)^{\dagger}\left(K-\sigma K_{G}\right)=\mathcal{P}_{\mathcal{R}\left(\left(K-\sigma K_{G}\right)^{T}\right)}=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)}, \tag{13}
\end{equation*}
$$

namely $\left(K-\sigma K_{G}\right)^{\dagger}\left(K-\sigma K_{G}\right)$ is an orthogonal projection onto $\mathcal{R}\left(\left(K-\sigma K_{G}\right)^{T}\right)=\mathcal{R}\left(K-\sigma K_{G}\right)$. Therefore, from (9) and (13),

$$
\left(K-\sigma K_{G}\right)^{\dagger}\left(K-\sigma K_{G}\right) W=W\left[\begin{array}{lll}
I_{n_{1}} & &  \tag{14}\\
& I_{n_{2}} & \\
& & 0
\end{array}\right] .
$$

From Equations (12) and (14), we have the eigenvalue decomposition (6) of $C$.
Lemma 2. The matrix C defined in (5) has the following properties:
(i) $\left(\lambda^{\#}, x\right)$ is an eigenpair of $K_{G}-\lambda^{\#} K$ with nonzero finite $\lambda^{\#}$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of $C$ with $\mu \neq 0$ and $\mu \neq 1$ and $x \in \mathcal{Z}_{\frac{\perp}{c}}^{\perp}$, where $\mu=\frac{1}{1-\sigma \alpha^{\dagger}}$.
(ii) $\left(\lambda^{\#}, x\right)$ is an eigenpair of $K_{G}-\lambda^{\#} K$ with $\lambda^{\#}=0$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of $C$ with $\mu=1$ and $x \in \mathcal{Z}_{c}^{\perp}$.
(iii) ( $\lambda^{\#}, x$ ) is an eigenpair of $K_{G}-\lambda^{\#} K$ with $\left|\lambda^{\#}\right|=\infty$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of $C$ with $\mu=0$ and $x \in \mathcal{Z}_{c}^{\perp}$.
(iv) If $x \in \mathcal{Z}_{c}, C x=0$.

Proof. The lemma can be proved by comparing the eigenvalue decomposition (6) of $C$ with the canonical form (3) of $K_{G}-\lambda^{\#} K$. Specifically, for (i) and (ii), recall that each column of $W_{1}$ is an eigenvector associated with a real, finite eigenvalue $\lambda^{\#}$ of the pencil $K_{G}-\lambda^{\#} K$ and the eigenvector is perpendicular to the common nullspace $\mathcal{Z}_{c}$. From (6), each column of $W_{1}$ is now an eigenvector associated with a non-zero, finite eigenvalue $\mu=\left(1-\sigma \lambda^{\#}\right)^{-1}$ of the eigenproblem (5).

To show (iii), recall that each column of $W_{2}$ is an eigenvector associated with an infinite eigenvalue of the pencil $K_{G}-\lambda^{\#} K$ and the eigenvector is perpendicular to the common nullspace $\mathcal{Z}_{c}$. From (6), each column of $W_{2}$ is now an eigenvector associated with zero eigenvalue of the eigenproblem (5).

Finally, for (iv), the common nullspace $\mathcal{Z}_{c}$ is spanned by the columns of $W_{3}$ and, from (6), we know that $C x=0$ if $x \in \mathcal{Z}_{c}$.

The following theorem provides the relationship of nontrivial eigenpairs between the original buckling eigenvalue problem (1) and the ordinary eigenvalue problem (5).

Theorem 2. $(\lambda, x)$ is an eigenpair of the pencil $K-\lambda K_{G}$ with nonzero finite eigenvalue $\lambda$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of the matrix $C$ in (5) with $\mu \neq 0$ and $\mu \neq 1$ and $x \in \mathcal{Z}_{c}^{\perp}$, where $\mu=\frac{\lambda}{\lambda-\sigma}$ and $\sigma \neq 0$.

Proof. Note that $(\lambda, x)$ is an eigenpair of $K-\lambda K_{G}$ with nonzero finite eigenvalue $\lambda$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $\left(\lambda^{\#}, x\right)$ is an eigenpair of $K_{G}-\lambda^{\#} K$ with non-zero finite eigenvalue $\lambda^{\#}=\lambda^{-1}$ and $x \in \mathcal{Z}_{c}^{\perp}$. Also, from Lemma 2 (i), we know that ( $\lambda^{\#}, x$ ) is an eigenpair of $K_{G}-\lambda^{\#} K$ with nonzero finite eigenvalue $\lambda^{\#}$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of the eigenvalue problem $C x=\mu x$ with $\mu=\frac{1}{1-\sigma \lambda^{\#}}, \mu \neq 0$ and $\mu \neq 1$, and $x \in \mathcal{Z}_{c}^{\perp}$. Therefore, $(\lambda, x)$ is an eigenpair of the pencil $K-\lambda K_{G}$ with nonzero finite eigenvalue $\lambda$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $(\mu, x)$ is an eigenpair of the eigenvalue problem $C x=\mu x$ with $\mu=\frac{\lambda}{\lambda-\sigma}, \mu \neq 0$ and $\mu \neq 1$, and $x \in \mathcal{Z}_{c}^{\perp}$.

By Theorem 2, near the shift $\sigma$, the eigenpairs $(\lambda, x)$ of $K-\lambda K_{G}$ with non-zero finite eigenvalues $\lambda$ and $x \in \mathcal{Z}_{c}^{\perp}$ are transformed into eigenpairs $(\mu, x)$ of $C$ with nonzero eigenvalues $\mu$, which typically are well-separated, and those away from the shift $\sigma$ are transformed into clustered eigenpairs $(\mu, x)$ of $C$ near unity as shown in Figure 1 . We note that the eigenpairs ( $\mu, x$ ) with $\mu=0$ or $\mu=1$ are not the ones of interest. The eigenpairs $(1, x)$ correspond to eigenpairs of $K-\lambda K_{G}$ with infinite eigenvalues and the eigenpairs $(0, x)$ correspond to eigenpairs of $K-\lambda K_{G}$ with $x \in \mathcal{N}(K)$.

## 2.3 | Regularization of the inner product

In this subsection we introduce a positive definite matrix $M$ from a low-rank updating of $K$, and then show that the matrix $C$ in the generalized buckling spectral transformation (5) is symmetric with respect to the inner product induced by $M$.

Theorem 3. Let $C$ be defined in (5). Let $Z=\left[Z_{N} Z_{C}\right]$ span the nullspace $\mathcal{N}(K)$ and $Z_{C}$ span the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$. Define

$$
\begin{equation*}
M=K+\left(K_{G} Z_{N}\right) H_{N}\left(K_{G} Z_{N}\right)^{T}+Z_{C} H_{C} Z_{C}^{T} \tag{15}
\end{equation*}
$$

where $H_{N}$ and $H_{C}$ are arbitrary positive definite matrices. Then
(i) the matrix $M$ is positive definite,
(ii) the matrix $C$ is symmetric with respect to the inner product induced by $M$.

Proof. By the canonical form (3), we have

FIGURE 1 Buckling spectral transformation with $\sigma<0$ (left) and $\sigma>0$ (right)



$$
\mathcal{N}(K)=\mathcal{R}\left(W_{2}\right) \oplus \mathcal{R}\left(W_{3}\right)=\mathcal{R}\left(Z_{N}\right) \oplus \mathcal{R}\left(Z_{C}\right) \quad \text { and } \quad \mathcal{Z}_{c}=\mathcal{R}\left(W_{3}\right)=\mathcal{R}\left(Z_{C}\right),
$$

and

$$
\left[\begin{array}{ll}
Z_{N} & Z_{C}
\end{array}\right]=\left[\begin{array}{ll}
W_{2} & W_{3}
\end{array}\right]\left[\begin{array}{cc}
R_{22} & O \\
R_{32} & R_{33}
\end{array}\right]
$$

for some matrices $R_{22} \in \mathbb{R}^{n_{2} \times n_{2}}, R_{32} \in \mathbb{R}^{n_{3} \times n_{2}}, R_{33} \in \mathbb{R}^{n_{3} \times n_{3}}$, and $R_{22}$ and $R_{33}$ are nonsingular. Therefore,

$$
W^{T} K_{G} Z_{N}=W^{T} K_{G}\left(W_{2} R_{22}+W_{3} R_{32}\right)=W^{T} K_{G} W_{2} R_{22}=\left[\begin{array}{c}
0 \\
\Lambda_{2}^{\#} R_{22} \\
0
\end{array}\right]
$$

Since the basis $W$ satisfies the condition (4),

$$
W^{T} Z_{C}=W^{T} W_{3} R_{33}=\left[\begin{array}{c}
0 \\
0 \\
\left(W_{3}^{T} W_{3}\right) R_{33}
\end{array}\right]
$$

Therefore,

$$
W^{T} M W=W^{T}\left(K+\left(K_{G} Z_{N}\right) H_{N}\left(K_{G} Z_{N}\right)^{T}+Z_{C} H_{C} Z_{C}^{T}\right) W=\left[\begin{array}{llll}
I_{n_{1}} & &  \tag{16}\\
& \widehat{H}_{N} & \\
& & \widehat{H}_{C}
\end{array}\right]
$$

where

$$
\widehat{H}_{N}=\Lambda_{2}^{\#} R_{22} H_{N} R_{22}^{T} \Lambda_{2}^{\#} \quad \text { and } \quad \hat{H}_{C}=\left(W_{3}^{T} W_{3}\right) R_{33} H_{C} R_{33}^{T}\left(W_{3}^{T} W_{3}\right)
$$

To prove that $M$ is positive definite, we show that both $\hat{H}_{N}$ and $\widehat{H}_{C}$ are positive definite. For the matrix $\hat{H}_{N}$, we note that the matrix $H_{N}$ is positive definite and the matrix $R_{22}$ is nonsingular. Also, from Theorem 1, the diagonal matrix $\Lambda_{2}^{\#}$ is nonsingular. Therefore, the matrix $\hat{H}_{N}$ is positive definite. For the matrix $\hat{H}_{C}$, we note that the matrix $H_{C}$ is positive definite and the matrix $R_{33}$ is nonsingular. Also, since the matrix $W_{3}$ is of full rank, the symmetric matrix $W_{3}^{T} W_{3}$ is nonsingular.

Therefore, the matrix $\hat{H}_{C}$ is also positive definite. This proves $(i)$.
To prove (ii), by the eigenvalue decomposition (6) of $C$ and (16), we have

$$
W^{T} M C W=W^{T} M W W^{-1} C W=\left[\begin{array}{lll}
\left(I_{n_{1}}-\sigma \Lambda_{1}^{\#}\right)^{-1} & & \\
& 0 & \\
& & 0
\end{array}\right] .
$$

Therefore, the matrix $M C$ is symmetric, which means that the matrix $C$ is symmetric with respect to the inner product induced by $M$.

Remark 2. We note that if the pencil $K-\lambda K_{G}$ is regular, Theorem 3 is still applicable. In this case, the matrix $C$ in (2) is symmetric with respect to the inner product induced by $M=K+\left(K_{G} Z_{N}\right) H_{N}\left(K_{G} Z_{N}\right)^{T}$.

## 3 | SHIFT-INVERT LANCZOS METHOD

## 3.1 | Shift-invert Lanczos method

Using Theorem 2, we have generalized the buckling spectral transformation to the singular pencil $K-\lambda K_{G}$ and converted the buckling eigenproblem (1) into an equivalent ordinary eigenvalue problem (5). From Theorem 3, we know that the
matrix $C$ in (5) is symmetric with respect to the inner product induced by the positive definite matrix $M$ in (15). It naturally leads that to solve the buckling eigenvalue problem (1), we can use the Lanczos method on the matrix $C$ with the inner product induced by $M$. This new strategy is also referred to as the shift-invert Lanczos method and outlined in Algorithm 1.

The shift-invert Lanczos method, after $j$ steps, computes a sequence of Lanczos vectors $\left\{v_{1} \ldots v_{j+1}\right\}$ and a symmetric tridiagonal matrix $T_{j}=\operatorname{tridiag}\left(\beta_{i-1}, \alpha_{i}, \beta_{i}\right)$ satisfying the governing equations

$$
\begin{equation*}
C V_{j}=V_{j} T_{j}+\beta_{j} v_{j+1} e_{j}^{T} \quad \text { and } \quad V_{j+1}^{T} M V_{j+1}=I_{j+1} \tag{17}
\end{equation*}
$$

where $V_{j+1} \equiv\left[\begin{array}{lll}v_{1} & \ldots & v_{j+1}\end{array}\right]$. Care must be taken to ensure that the equations in (17) are satisfied ${ }^{7-9,14,15}$ in the presence of finite-precision arithmetic. In particular, at step 11 of Algorithm 1, we perform full re-orthogonalization at each iteration using the classical Gram-Schmidt process (Reference 16, p. 120), that is,

$$
r=r-V_{j}\left(V_{j}^{T}(M r)\right)
$$

Efficient practical techniques such as partial and selective re-orthogonalization have been developed and well-implemented. ${ }^{9,14,15}$ In next subsection, we will focus on the implementation of the matrix-vector product $u=C v$ at step 6.

```
Algorithm 1. Shift-invert Lanczos method for the buckling eigenvalue problem (1)
\(\left(\widehat{\mu}_{i}, \widehat{x}_{i}\right)\) of \(C\).
    \(r=v\)
    \(p=M r\)
    \(\beta_{0}=\left(p^{T} r\right)^{1 / 2}\)
    for \(j=1,2, \ldots\) do
    \(v_{j}=r / \beta_{j-1}\)
    \(r=C v_{j}\) (see Section 3.2)
    \(r=r-\beta_{j-1} v_{j-1}\)
    \(p=M r\)
    \(\alpha_{j}=v_{j}^{T} p\)
    \(r=r-\alpha_{j} v_{j}\)
    perform re-orthogonalization if necessary
    \(p=M r\)
    \(\beta_{j}=\left(p^{T} r\right)^{1 / 2}\)
    compute the eigenpairs \(\left(\widehat{\mu}_{i}, \widehat{s}_{i}\right)\) of \(T_{j}=\operatorname{tridiag}\left(\beta_{i-1}, \alpha_{i}, \beta_{i}\right)\)
    use \(t o l\) to check the relative residual norm (33) for convergence
    end for
    return the converged eigenpairs \(\left(\mu_{i}, \widehat{x}_{i}=V_{j} \widehat{s}_{i}\right)\)
```

This algorithm takes as input the starting vector $v$, the matrix-vector product $u=C v$, the matrix $M$ with the positive def-
inite matrices $H_{N}$ and $H_{C}$, and the tolerance value tol for the relative residual norm. It returns the converged eigenpairs

## 3.2 | The matrix-vector product

We first show that the matrix-vector product $u=C v=\left(K-\sigma K_{G}\right)^{\dagger} K v$ is connected with the solution of a consistent singular linear system with constraint.

Theorem 4. Given $v \in \mathbb{R}^{n}$, the vector

$$
\begin{equation*}
u=\left(K-\sigma K_{G}\right)^{\dagger} K v, \tag{18}
\end{equation*}
$$

is the unique solution of the consistent singular linear system

$$
\begin{equation*}
\left(K-\sigma K_{G}\right) u=K v \tag{19}
\end{equation*}
$$

with the constraint

$$
\begin{equation*}
Z_{C}^{T} u=0 \tag{20}
\end{equation*}
$$

where $Z_{C}$ is a basis of the common nullspace of $K$ and $K_{G}$.

Proof. First note that since both $K$ and $K-\sigma K_{G}$ are symmetric, we have

$$
\begin{equation*}
\mathcal{R}(K)=\mathcal{N}(K)^{\perp} \quad \text { and } \quad \mathcal{R}\left(K-\sigma K_{G}\right)=\mathcal{N}\left(K-\sigma K_{G}\right)^{\perp}=\mathcal{Z}_{c}^{\perp} \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathcal{Z}_{c}=\mathcal{N}\left(K-\sigma K_{G}\right) \subset \mathcal{N}(K) \tag{22}
\end{equation*}
$$

Therefore from (21) and (22),

$$
K v \in \mathcal{R}(K) \subset \mathcal{R}\left(K-\sigma K_{G}\right)
$$

which implies that the linear system (19) is consistent. From (18),

$$
\begin{equation*}
\left(K-\sigma K_{G}\right) u=\left(K-\sigma K_{G}\right)\left(K-\sigma K_{G}\right)^{\dagger} K v=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)} K v=K v, \tag{23}
\end{equation*}
$$

where $\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)}$ is an orthogonal projection onto $\mathcal{R}\left(K-\sigma K_{G}\right)$ (by the Moore-Penrose conditions Reference 13, p. 290). This means that $u$ is a solution of the consistent singular linear system (19).

On the other hand, from (18) and (23),

$$
\begin{equation*}
u=\left(K-\sigma K_{G}\right)^{\dagger} K v=\left(K-\sigma K_{G}\right)^{\dagger}\left(K-\sigma K_{G}\right) u=\mathcal{P}_{\mathcal{R}\left(\left(K-\sigma K_{G}\right)^{T}\right)} u=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)} u \tag{24}
\end{equation*}
$$

Since $\mathcal{R}\left(K-\sigma K_{G}\right)=\mathcal{Z}_{c}^{\perp}$, it implies that $u$ is perpendicular to the common nullspace $\mathcal{Z}_{c}$, which is also the nullspace $\mathcal{N}\left(K-\sigma K_{G}\right)$.

The uniqueness can be shown as follows. Given two solutions $u_{1}$ and $u_{2}$ to (19), the difference $u_{1}-u_{2}$ would satisfy $\left(K-\sigma K_{G}\right)\left(u_{1}-u_{2}\right)=0$, which implies $u_{1}-u_{2} \in \mathcal{Z}_{c}$. However, since both solutions satisfy the constraint (20), $Z_{C}^{T}\left(u_{1}-\right.$ $\left.u_{2}\right)=0$. Therefore, $u_{1}-u_{2}=0$.

We now present method to compute the matrix-vector product $u=C v$. First, we have the following theorem to extract a non-singular submatrix of $K-\sigma K_{G}$ by exploiting the basis $Z_{C}$.

Theorem 5. Let $Z_{C} \in \mathbb{R}^{n \times n_{3}}$ be a basis of $\mathcal{N}\left(K-\sigma K_{G}\right)$ and $P \in \mathbb{R}^{n \times n}$ be a permutation matrix such that $P^{T} Z_{C} \equiv\left[\begin{array}{l}Y_{1} \\ Y_{2}\end{array}\right]$, and $Y_{2} \in \mathbb{R}^{n_{3} \times n_{3}}$ is nonsingular. Define

$$
S=P^{T}\left(K-\sigma K_{G}\right) P \quad \text { and } \quad S={ }_{n_{3}}^{n-n_{3}}\left[\begin{array}{cc}
n-n_{3} & n_{3}  \tag{25}\\
S_{11}^{\sigma} & S_{12} \\
S_{12}^{T} & S_{22}
\end{array}\right] .
$$

Then
(1) the submatrix $S_{11}^{\sigma} \in \mathbb{R}^{\left(n-n_{3}\right) \times\left(n-n_{3}\right)}$ is nonsingular,
(2) $v_{+}\left(S_{11}^{\sigma}\right)=v_{+}\left(K-\sigma K_{G}\right)$ and $v_{-}\left(S_{11}^{\sigma}\right)=v_{-}\left(K-\sigma K_{G}\right)$, where $v_{+}(X)$ and $v_{-}(X)$ denote the numbers of positive and negative eigenvalues of the symmetric matrix $X$, respectively.

Proof. Let

$$
\left.E=\begin{array}{c}
n-n_{3} \\
n_{3}
\end{array} \begin{array}{cc}
n-n_{3} & n_{3} \\
I_{n-n_{3}} & Y_{1} \\
0 & Y_{2}
\end{array}\right] \in \mathbb{R}^{n \times n} .
$$

The matrix $E$ is non-singular since $Y_{2}$ is nonsingular. By the congruence transformation, we have

$$
E^{T} S E=E^{T} P^{T}\left(K-\sigma K_{G}\right) P E=E^{T}\left[\begin{array}{cc}
S_{11}^{\sigma} & S_{12}  \tag{26}\\
S_{12}^{T} & S_{22}
\end{array}\right] E={ }_{{ }_{n}-n_{3}}^{n_{3}}\left[\begin{array}{cc}
n-n_{3} & n_{3} \\
S_{11}^{\sigma} & 0 \\
0 & 0
\end{array}\right] .
$$

Sylvester's law (Reference 13, p. 448) tells that the matrices $K-\sigma K_{G}$ and $E^{T} S E$ have the same inertias. In particular, from (26), we know that

$$
v_{+}\left(K-\sigma K_{G}\right)=v_{+}\left(S_{11}^{\sigma}\right), \quad v_{-}\left(K-\sigma K_{G}\right)=v_{-}\left(S_{11}^{\sigma}\right),
$$

and

$$
\begin{equation*}
v_{0}\left(K-\sigma K_{G}\right)=v_{0}\left(S_{11}^{\sigma}\right)+n_{3} \tag{27}
\end{equation*}
$$

But $v_{0}\left(K-\sigma K_{G}\right)=\operatorname{dim}\left(\mathcal{N}\left(K-\sigma K_{G}\right)\right)=n_{3}$. Therefore, from (27), $v_{0}\left(S_{11}^{\sigma}\right)=0$ and $S_{11}^{\sigma}$ is nonsingular.
Theorem 5 was inspired by Reference 17, theorem 2.2, where the authors consider solving a consistent semi-definite linear systems $A x=b$ from the electromagnetic applications. ${ }^{18}$ The matrix $A$, generated from the finite element modeling, is positive semi-definite and an explicit basis of the nullspace of $A$ is available. This explicit basis of the nullspace is then used to identify a nonsingular part of $A$ and a solution of the linear system can be computed from it. Although in the buckling eigenvalue probem (1), the matrix $K-\sigma K_{G}$ is indefinite, we found that the strategy developed in Reference 17 can be generalized to the system (19) and (20).

By Theorem 5, the method to solve (19), that is, compute the matrix-vector product $u=C v=\left(K-\sigma K_{G}\right)^{\dagger} K v$, can be described in two steps:

1. Find a solution $u_{p}$ of the consistent singular linear system (19).
2. Compute $u=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)} u_{p}$ to satisfy the constraint (20), where $\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)}$ is an orthogonal projection onto $\mathcal{R}\left(K-\sigma K_{G}\right)$.

Specifically, in Step 1, find the permutation matrix $P$ as described in Theorem 5, and rewrite (19) in the partitioned form (25):

$$
\left[\begin{array}{ll}
S_{11}^{\sigma} & S_{12}  \tag{28}\\
S_{12}^{T} & S_{22}
\end{array}\right]\left[\begin{array}{l}
w_{1} \\
w_{2}
\end{array}\right]=\left[\begin{array}{l}
c_{1} \\
c_{2}
\end{array}\right] \in \mathcal{R}(S),
$$

where

$$
\left[\begin{array}{l}
w_{1} \\
w_{2}
\end{array}\right] \equiv P^{T} u \quad \text { and } \quad\left[\begin{array}{l}
c_{1} \\
c_{2}
\end{array}\right] \equiv P^{T} K v .
$$

Since $S_{11}^{\sigma}$ is nonsingular, $S_{11}^{\sigma}$ is of full rank and the leading $n-n_{3}$ columns of $S$ are linearly independent. On the other hand, we know that $\operatorname{rank}(S)=\operatorname{rank}\left(K-\sigma K_{G}\right)=n-n_{3}$. Therefore, the leading $n-n_{3}$ columns of $S$ is a basis of $\mathcal{R}(S)$, and there is a solution $w_{p}$ of (28) with $w_{2}=0$. Direct substitution gives

$$
w_{p}=\left[\begin{array}{c}
\left(S_{11}^{\sigma}\right)^{-1} c_{1} \\
0
\end{array}\right],
$$

where the inverse $\left(S_{11}^{\sigma}\right)^{-1}$ can be computed using the sparse $\operatorname{LDL}^{\mathrm{T}}$ factorization of $S_{11}^{\sigma}{ }^{19,20}$ A solution $u_{p}$ of (19) is then given by

$$
u_{p}=P\left[\begin{array}{c}
\left(S_{11}^{\sigma}\right)^{-1} c_{1} \\
0
\end{array}\right]
$$

In Step 2 , since $Z_{C}$ is a basis of $\mathcal{N}\left(K-\sigma K_{G}\right)$, which is the orthogonal complement to $\mathcal{R}\left(K-\sigma K_{G}\right)$, the vector $u$ can be computed by the projection

$$
u=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)} u_{p}=\left(I-Z_{C}\left(Z_{C}^{T} Z_{C}\right)^{-1} Z_{C}^{T}\right) u_{p}
$$

If $Z_{C}$ is an orthonormal basis, then

$$
u=\mathcal{P}_{\mathcal{R}\left(K-\sigma K_{G}\right)} u_{p}=\left(I-Z_{C} Z_{C}^{T}\right) u_{p}
$$

## 4 | COUNTING EIGENVALUES

In this section, as a validation scheme, we discuss a way to count the number of eigenvalues in a given interval. In the following, $v_{+}(A)$ and $v_{-}(A)$ denote the number of positive and negative eigenvalues of a symmetric matrix $A$, respectively. $n(\alpha, \beta)$ and $n^{\#}(\alpha, \beta)$ denote the numbers of eigenvalues of the pencil $K-\lambda K_{G}$ and the reversed pencil $K_{G}-\lambda^{\#} K$ in an interval $(\alpha, \beta)$, respectively.

First, we consider the following lemma.
Lemma 3. Let $Z=\left[Z_{N} Z_{C}\right]$ span the nullspace $\mathcal{N}(K)$ and $Z_{C}$ span the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$, then
(i) for $\alpha<0, n(\alpha, 0)=v_{-}\left(K-\alpha K_{G}\right)-v_{-}\left(Z_{N}^{T} K_{G} Z_{N}\right)$,
(ii) for $\alpha>0, n(0, \alpha)=v_{-}\left(K-\alpha K_{G}\right)-v_{+}\left(Z_{N}^{T} K_{G} Z_{N}\right)$.

In addition, the matrix $Z_{N}^{T} K_{G} Z_{N}$ is nonsingular.

Proof. The proof is based on the following two facts: (1) $(\lambda, x)$ is an eigenpair of the pencil $K-\lambda K_{G}$ with nonzero finite eigenvalue $\lambda$ and $x \in \mathcal{Z}_{c}^{\perp}$ if and only if $\left(\lambda^{\#}, x\right)$ is an eigenpair of the pencil $K_{G}-\lambda^{\#} K$ with nonzero finite eigenvalue $\lambda^{\#}=\lambda^{-1}$ and $x \in \mathcal{Z}_{c}^{\perp}$. (2) By the canonical form (3), we have

$$
W^{T}\left(K_{G}-\frac{1}{\alpha} K\right) W=\left[\begin{array}{llll}
\Lambda_{1}^{\#}-\frac{1}{\alpha} I_{n_{1}} & & \\
& \Lambda_{2}^{\#} & \\
& & 0
\end{array}\right]
$$

Consequently, by Sylvester's law, we have

$$
\begin{aligned}
& v_{-}\left(K_{G}-\frac{1}{\alpha} K\right)=v_{-}\left(\Lambda_{1}^{\#}-\frac{1}{\alpha} I_{n_{1}}\right)+v_{-}\left(\Lambda_{2}^{\#}\right) \\
& v_{+}\left(K_{G}-\frac{1}{\alpha} K\right)=v_{+}\left(\Lambda_{1}^{\#}-\frac{1}{\alpha} I_{n_{1}}\right)+v_{+}\left(\Lambda_{2}^{\#}\right)
\end{aligned}
$$

Now, for (i), since $\alpha<0$,

$$
\begin{equation*}
n(\alpha, 0)=n^{\#}\left(-\infty, \frac{1}{\alpha}\right)=v_{-}\left(\Lambda_{1}^{\#}-\frac{1}{\alpha} I_{n_{1}}\right)=v_{-}\left(K_{G}-\frac{1}{\alpha} K\right)-v_{-}\left(\Lambda_{2}^{\#}\right)=v_{-}\left(K-\alpha K_{G}\right)-v_{-}\left(\Lambda_{2}^{\#}\right) \tag{29}
\end{equation*}
$$

where for the second equality, see Remark 1. For (ii), since $\alpha>0$,

$$
\begin{equation*}
n(0, \alpha)=n^{\#}\left(\frac{1}{\alpha},+\infty\right)=v_{+}\left(\Lambda_{1}^{\#}-\frac{1}{\alpha} I_{n_{1}}\right)=v_{+}\left(K_{G}-\frac{1}{\alpha} K\right)-v_{+}\left(\Lambda_{2}^{\#}\right)=v_{-}\left(K-\alpha K_{G}\right)-v_{+}\left(\Lambda_{2}^{\#}\right) . \tag{30}
\end{equation*}
$$

On the other hand, by the canonical form (3), we have

$$
\mathcal{N}(K)=\mathcal{R}\left(Z_{N}\right) \oplus \mathcal{R}\left(Z_{C}\right)=\mathcal{R}\left(W_{2}\right) \oplus \mathcal{R}\left(W_{3}\right) \quad \text { and } \quad \mathcal{Z}_{c}=\mathcal{R}\left(Z_{C}\right)=\mathcal{R}\left(W_{3}\right)
$$

and

$$
Z_{N}=W_{2} R_{22}+W_{3} R_{32}
$$

where $R_{22} \in \mathbb{R}^{n_{2} \times n_{2}}, R_{32} \in \mathbb{R}^{n_{3} \times n_{2}}$ and $R_{22}$ is nonsingular. Also, we know that $W_{2}^{T} K_{G} W_{2}=\Lambda_{2}^{\#}$. Therefore,

$$
Z_{N}^{T} K_{G} Z_{N}=R_{22}^{T}\left(W_{2}^{T} K_{G} W_{2}\right) R_{22}=R_{22}^{T} \Lambda_{2}^{\#} R_{22}
$$

This implies that the matrix $Z_{N}^{T} K_{G} Z_{N}$ is nonsingular, and by Sylvester's law, we have

$$
\begin{equation*}
v_{-}\left(\Lambda_{2}^{\#}\right)=v_{-}\left(Z_{N}^{T} K_{G} Z_{N}\right) \quad \text { and } \quad v_{+}\left(\Lambda_{2}^{\#}\right)=v_{+}\left(Z_{N}^{T} K_{G} Z_{N}\right) \tag{31}
\end{equation*}
$$

The lemma is an immediate consequence of (29), (30), and (31).

Lemma 3 establishes the relation between the number of eigenvalues in the interval $(\alpha, 0)$ or $(0, \alpha)$ and the inertia $v_{-}\left(K-\alpha K_{G}\right)$. Below, we discuss how to express the inertia $v_{-}\left(K-\alpha K_{G}\right)$ in terms of the submatrix $S_{11}^{\alpha}$ in (25).

Lemma 4. In terms of the submatrix $S_{11}^{\alpha}$ in (25),

$$
\begin{equation*}
v_{-}\left(K-\alpha K_{G}\right)=v_{-}\left(S_{11}^{\alpha}\right) \tag{32}
\end{equation*}
$$

Proof. The equality (32) immediately follows from Theorem 5.

Combining Lemmas 3 and 4, we have the following theorem which provides a computational approach to count the number of eigenvalues of $K-\lambda K_{G}$ using the inertias of $S_{11}^{\alpha}$.

Theorem 6. In terms of the submatrix $S_{11}^{\alpha}$ in (25), we have
(i) $n(\alpha, 0)=v_{-}\left(S_{11}^{\alpha}\right)-v_{-}\left(Z_{N}^{T} K_{G} Z_{N}\right)$, if $\alpha<0$.
(ii) $n(0, \alpha)=\nu_{-}\left(S_{11}^{\sigma}\right)-\nu_{+}\left(Z_{N}^{T} K_{G} Z_{N}\right)$, if $\alpha>0$.

Remark 3. In practice, the inertia $v_{-}\left(S_{11}^{\alpha}\right)$ is a by-product of the sparse LDL $^{\mathrm{T}}$ factorizations of the submatrix $S_{11}^{\alpha}$ (Reference 21, p. 214). The inertias $v_{-}\left(Z_{N}^{T} K_{G} Z_{N}\right)$ and $v_{+}\left(Z_{N}^{T} K_{G} Z_{N}\right)$ can be easily computed since the size of $Z_{N}^{T} K_{G} Z_{N}$ is typically small in buckling analysis.

## 5 | NUMERICAL EXAMPLES

In this section, we begin with a synthetic example to illustrate the issue associated with the growth of the norms of the Lanczos vectors with $K$-inner product and the consequence of the growth as discussed by Meerbergen ${ }^{4}$ and Stewart. ${ }^{5}$ Then we demonstrate the efficacy of the proposed shift-invert Lanczos method for an example arising in industrial buckling analysis of structures.


FIGURE 2 Left: the 2-norms of the Lanczos vectors $v_{j}$. Middle: the relative residual norms of the approximate eigenpairs $\left(\hat{\lambda}_{i}, \hat{x}_{i}\right)$. Right: the 2-norms of the Lanczos vectors $v_{j}$ with ( + ) and without ( x ) implicit restart

Algorithm 1 is implemented in MATLAB. ${ }^{\text {a }}$ The accuracy of a computed eigenpair $\left(\hat{\lambda}_{i}, \widehat{x}_{i}\right)$ of the buckling eigenvalue problem (1) is measured by the relative residual norm

$$
\begin{equation*}
\eta\left(\hat{\lambda}_{i}, \widehat{x}_{i}\right) \equiv \frac{\left\|K \widehat{x}_{i}-\widehat{\lambda}_{i} K_{G} \widehat{x}_{i}\right\|_{2}}{\left(\|K\|_{1}+\left|\hat{\lambda_{i}}\right|\left\|K_{G}\right\|_{1}\right)\left\|\hat{x}_{i}\right\|_{2}} . \tag{33}
\end{equation*}
$$

The Euclidean angle $\theta_{i}=\angle\left(\widehat{x}_{i}, \mathcal{Z}_{c}\right)$ is computed for checking if $\widehat{x}_{i}$ is perpendicular to the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G} .{ }^{22,23}$

Example 1. Let us consider the following matrix pair ( $K, K_{G}$ ) similar to the ones constructed by Meerbergen ${ }^{4}$ and Stewart: ${ }^{5}$

$$
K=Q \Lambda Q^{T} \in \mathbb{R}^{n \times n} \quad \text { and } \quad K_{G}=Q \Phi Q^{T} \in \mathbb{R}^{n \times n}
$$

where $Q \in \mathbb{R}^{n \times n}$ is a random orthogonal matrix, $\Lambda \in \mathbb{R}^{n \times n}$ and $\Phi \in \mathbb{R}^{n \times n}$ are diagonal matrices with diagonal elements

$$
\Lambda_{k k}=\left\{\begin{array}{ll}
k, & \text { if } 1 \leq k \leq n-m \\
0, & \text { otherwise }
\end{array} \quad \text { and } \quad \Phi_{k k}=(-1)^{k}, \quad 1 \leq k \leq n\right.
$$

By construction, $K$ is positive semi-definite and $K_{G}$ is indefinite, and the pencil $K-\lambda K_{G}$ is regular. The last $m$ columns of $Q$ form a basis of the nullspace $\mathcal{N}(K)$. For $1 \leq k \leq n-m$, the $k$ th column of $Q$ is an eigenvector and the associated eigenvalue is $\lambda_{k}=(-1)^{k} \cdot k$. The zero eigenvalue of $C \equiv\left(K-\sigma K_{G}\right)^{-1} K$ is a well-separated eigenvalue, and the associated eigenspace is also the nullspace of $K$. We use the MATLAB function $1 d 1$ to compute the LDL $^{\mathrm{T}}$ factorization of the shifted matrix $K-\sigma K_{G}$.

For numerical experiments, we take $n=500$ and $m=1$. We use the buckling spectral transformation (2) with the shift $\sigma=-0.6$. We run the Lanczos method with $K$-inner product, and the starting vector $C x_{0}$ with $x_{0}=[1, \ldots, 1]^{T}$. The approximate eigenpairs $\left(\hat{\lambda}_{i}, \widehat{x}_{i}\right)$ of (1) are computed by $\left(\hat{\lambda}_{i}, \widehat{x}_{i}\right)=\left(\frac{\sigma \hat{\mu}_{i}}{\hat{\mu}_{i}-1}, \widehat{x}_{i}\right)$.

The left plot of Figure 2 shows the 2-norms of 40 Lanczos vectors $v_{j}$. As observed by Meerbergen ${ }^{4}$ and Stewart, ${ }^{5}$ the 2 -norms of Lanczos vectors $v_{j}$ grow rapidly. Consequently, as shown in the middle plot of Figure 2, the accuracy of approximate eigenpairs $\left(\hat{\lambda}_{i}, \widehat{x}_{i}\right)$ deteriorates. In contrast, when we replace the $K$-inner product by the positive definite $M$-inner product with $H_{N}=I_{m}$, we observe that the 2-norms of the Lanczos vectors are well bounded. Multiple eigenvalues near the shift $\sigma$ are computed with the relative residual norms around the machine precision.

We note that in Reference 4, Meerbergen proposed to control the norms of the Lanczos vectors by applying implicit restart. We experimented the schemes with and without the implicit restart The results are shown in the right plot of Figure 2 . We can see that the 2-norms of the Lanczos vectors still grow rapidly.
${ }^{\mathrm{a}} \mathrm{An}$ implementation is available at https://github.com/cplin722/bucklingEigs.

Example 2. This is an example from the buckling analysis of a finite element model of an airplane shown in Figure 3. The size of the pencil $K-\lambda K_{G}$ is $n=67,512$. The stiffness matrix $K$ is positive semi-definite and the dimension of the nullspace $\mathcal{N}(K)$ is known to be 6 , which corresponds to the six rigid body modes. ${ }^{2}$ The geometric stiffness matrix $K_{G}$ is symmetric but indefinite. The basis $Z$ of $\mathcal{N}(K)$ is computed by $Z=\left[-\left(K_{11}^{-1} K_{12}\right)^{T} I_{6}\right]^{T}$, ${ }^{2}$ where $\left[K_{11} K_{12}\right] \in \mathbb{R}^{(n-6) \times n}$ is the leading block rows of $K$. The dimension of the common nullspace $\mathcal{Z}_{c}$ of $K$ and $K_{G}$ is 3 , which can be easily computed from the basis $Z$, see Reference 13, theorem 6.4.1). The accuracy of the bases is shown in the table in Figure 3. We are interested in computing the nonzero eigenvalues of the pencil $K-\lambda K_{G}$ in an interval around zero and the associated eigenvectors perpendicular to the common nullspace $\mathcal{Z}_{c}$.

We use the method to compute the matrix-vector product $u=C v$ described in Section 3.2. We determine the permutation matrix $P$ by maximizing the number of nonzero entries in the last $n_{3}$ columns of $S$ in (25). The MATLAB function ldl, which uses MA57 ${ }^{24}$ for real sparse matrices, is used to compute the sparse LDL $^{\mathrm{T}}$ factorization of the submatrix $S_{11}^{\sigma}$. The pivot tolerance $\tau=0.1$ is used to control the numerical stability of the factorization. ${ }^{24}$ In defining the positive definite matrix $M$, we form the product $K_{G} Z_{N}$ and normalize each column of the matrices $K_{G} Z_{N}$ and $Z_{C}$. The condition number of $K_{G} Z_{N}$ after the normalization is $\kappa_{2}\left(K_{G} Z_{N}\right)=1.03$. Then we set the matrices $H_{N}=\omega I_{n_{2}}$ and $H_{C}=\omega I_{n_{3}}, \omega=\|K\|_{1}$, to balance the matrix $M .{ }^{25}$ The starting vector of the Lanczos procedure is $v=C x_{0}$ with $x_{0}$ being a random vector. ${ }^{8}$

To monitor the progress of the shift-invert Lanczos method, an approximate eigenpair ( $\hat{\mu}_{i}, \hat{x}_{i}$ ) computed from an eigenpair ( $\hat{\mu}_{i}, \hat{S}_{i}$ ) of the reduced matrix $T_{j}$ is considered to have converged if the following two conditions are satisfied:

$$
\begin{equation*}
\left|\widehat{\mu}_{i}\right| \geq \text { tol } \quad \text { and } \quad \frac{|\sigma|}{\left(\hat{\mu}_{i}-1\right)^{2}}\left|\beta_{j}\right|\left|e_{j}^{T} \widehat{s}_{i}\right|<\text { tol }, \tag{34}
\end{equation*}
$$

where the first condition excludes the zero eigenvalues and the second condition bounds the error of the computed eigenvalue $\hat{\lambda}_{i}=\frac{\sigma \hat{\mu}_{i}}{\hat{\mu}_{i}-1}$ with the prescribed tolerance $t o l$ (see References 7,9 , and 26, p. 357). In this numerical example, we experiment with the tolerance tol $=10^{-6}$.

We now show the numerical results for computing nonzero eigenvalues of the pencil $K-\lambda K_{G}$ and corresponding eigenvectors perpendicular to the common nullspace $\mathcal{Z}_{c}$ in the interval $(-8,8)$. First, let us consider the left-half interval $(-8,0)$. With the shift $\sigma=-4.0$, the shift-invert Lanczos method (Algorithm 1) computed 12 eigenvalues to the machine precision in the interval $(-8,0)$ at 38 th iteration. The accuracy of the computed eigenpairs $\left(\widehat{\lambda}_{i}=\frac{\sigma \widehat{\mu}_{i}}{\hat{\mu}_{i}-1}, \widehat{x}_{i}\right)$ is shown in Table 1. To validate the number of eigenvalues in the interval $(-8,0)$, we use the counting scheme described in Section 4. Using the inertias of the submatrix $S_{11}^{\alpha}$ with $\alpha=-8$ and Theorem 5, we have

$$
n(-8,0)=v_{-}\left(S_{11}^{\alpha}\right)-v_{-}\left(Z_{N}^{T} K_{G} Z_{N}\right)=15-3=12 .
$$

This matches the number of eigenvalues found in the interval.
Next let us consider the right-half interval ( 0,8 ). In this case, we use the shift $\sigma=4.0$. By the shift-invert Lanczos method (Algorithm 1), we found 13 eigenvalues to the machine precision in the interval $(0,8)$ at 44th iteration. The accuracy of the computed eigenpairs $\left(\widehat{\lambda}_{i}=\frac{\sigma \hat{\mu}_{i}}{\widehat{\mu}_{i}-1} \widehat{x}_{i}\right)$ are shown in Table 2 . To validate the number of eigenvalues in the interval $(0,8)$, we again use the counting scheme described in Section 4. Using the inertias of the submatrix $S_{11}^{\alpha}$ with $\alpha=8$ and Theorem 5, we have


| $i$ | $d_{i} /\left\\|K_{G}\right\\|_{1}$ | $\frac{\left\\|K z_{i}\right\\|_{2}}{\\|K\\|_{1}\left\\|_{i}\right\\|_{2}}$ | $\frac{\left\\|K_{G} z_{2}\right\\|_{2}}{\left\\|K_{G}\right\\|\left\\|_{1}\right\\| z_{i} \\|_{2}}$ |
| :---: | :---: | :---: | :---: |
| 1 | $6.90 \cdot 10^{-5}$ | $2.74 \cdot 10^{-16}$ | $6.78 \cdot 10^{-5}$ |
| 2 | $3.25 \cdot 10^{-5}$ | $4.88 \cdot 0^{-16}$ | $9.06 \cdot 10^{-6}$ |
| 3 | $2.32 \cdot 10^{-5}$ | $4.71 \cdot 10^{-16}$ | $1.19 \cdot 10^{-5}$ |
| 4 | $7.32 \cdot 10^{-16}$ | $2.68 \cdot 10^{-17}$ | $5.01 \cdot 10^{-18}$ |
| 5 | $1.26 \cdot 10^{-16}$ | $1.90 \cdot 10^{-17}$ | $4.89 \cdot 10^{-18}$ |
| 6 | $7.81 \cdot 10^{-18}$ | $2.37 \cdot 10^{-17}$ | $5.00 \cdot 10^{-18}$ |

FIGURE 3 Left: Finite element model of an airplane. Right: Accuracy of the bases for the nullspace of $K$ and common nullspace of $K$ and $K_{G}$. The second column shows the singular values $d_{i}$ of $K_{G} Y$ with $Y$ being an orthonormal basis of $\mathcal{N}(K)$. The third and fourth columns show the accuracy of the basis $Z=\left[\begin{array}{lll}Z_{N} & Z_{C}\end{array}\right]=\left[\begin{array}{llll}z_{1} & z_{2} & \ldots & z_{6}\end{array}\right]$

| $\boldsymbol{i}$ | $\hat{\lambda}_{\boldsymbol{i}}$ | $\boldsymbol{\eta}\left(\hat{\lambda}_{i}, \widehat{\boldsymbol{x}}_{\boldsymbol{i}}\right)$ | $\boldsymbol{\operatorname { c o s }} \angle\left(\hat{\boldsymbol{x}}_{\boldsymbol{i}}, \mathcal{Z}_{\boldsymbol{c}}\right)$ |
| ---: | :--- | :--- | :--- |
| 1 | -2.716598 | $1.48 \cdot 10^{-17}$ | $8.52 \cdot 10^{-17}$ |
| 2 | -2.883589 | $1.73 \cdot 10^{-17}$ | $8.27 \cdot 10^{-17}$ |
| 3 | -3.292700 | $1.37 \cdot 10^{-17}$ | $4.84 \cdot 10^{-18}$ |
| 4 | -3.378406 | $1.01 \cdot 10^{-17}$ | $2.38 \cdot 10^{-17}$ |
| 5 | -5.754628 | $2.72 \cdot 10^{-17}$ | $4.04 \cdot 10^{-17}$ |
| 6 | -5.854071 | $2.92 \cdot 10^{-17}$ | $3.47 \cdot 10^{-17}$ |
| 7 | -6.089281 | $3.14 \cdot 10^{-17}$ | $2.47 \cdot 10^{-17}$ |
| 8 | -6.228974 | $2.67 \cdot 10^{-17}$ | $6.24 \cdot 10^{-17}$ |
| 9 | -6.784766 | $5.33 \cdot 10^{-16}$ | $4.93 \cdot 10^{-17}$ |
| 10 | -6.886759 | $2.57 \cdot 10^{-15}$ | $7.67 \cdot 10^{-18}$ |
| 11 | -7.561377 | $1.88 \cdot 10^{-12}$ | $1.31 \cdot 10^{-16}$ |
| 12 | -7.745144 | $3.83 \cdot 10^{-12}$ | $4.87 \cdot 10^{-17}$ |


| $\boldsymbol{i}$ | $\hat{\lambda}_{\boldsymbol{i}}$ | $\boldsymbol{\eta}\left(\hat{\lambda}_{\boldsymbol{i}}, \hat{\boldsymbol{x}}_{\boldsymbol{i}}\right)$ | $\boldsymbol{\operatorname { c o s }} \angle\left(\hat{\boldsymbol{x}}_{\boldsymbol{i}}, \mathcal{Z}_{\boldsymbol{c}}\right)$ |
| ---: | :--- | :--- | :--- |
| 1 | 2.967043 | $3.80 \cdot 10^{-17}$ | $1.10 \cdot 10^{-16}$ |
| 2 | 3.025965 | $2.96 \cdot 10^{-17}$ | $3.39 \cdot 10^{-17}$ |
| 3 | 3.917831 | $1.71 \cdot 10^{-17}$ | $7.71 \cdot 10^{-17}$ |
| 4 | 4.008941 | $1.61 \cdot 10^{-17}$ | $7.13 \cdot 10^{-17}$ |
| 5 | 4.591063 | $2.43 \cdot 10^{-17}$ | $4.29 \cdot 10^{-17}$ |
| 6 | 4.662575 | $2.64 \cdot 10^{-17}$ | $2.47 \cdot 10^{-17}$ |
| 7 | 5.699271 | $5.24 \cdot 10^{-17}$ | $7.45 \cdot 10^{-17}$ |
| 8 | 5.725937 | $7.44 \cdot 10^{-17}$ | $1.38 \cdot 10^{-17}$ |
| 9 | 6.465175 | $7.40 \cdot 10^{-16}$ | $1.14 \cdot 10^{-16}$ |
| 10 | 6.598173 | $7.96 \cdot 10^{-15}$ | $2.18 \cdot 10^{-16}$ |
| 11 | 7.285975 | $4.45 \cdot 10^{-15}$ | $3.32 \cdot 10^{-16}$ |
| 12 | 7.626265 | $2.41 \cdot 10^{-14}$ | $1.39 \cdot 10^{-15}$ |
| 13 | 7.880296 | $1.24 \cdot 10^{-12}$ | $3.71 \cdot 10^{-14}$ |

TABLE 2 Results of 13 computed eigenvalues in the interval $(0,8)$ after 44 steps of the Lanczos method with the shift $\sigma=4.0$. $\left\|\widehat{X}^{T} M \widehat{X}-I_{13}\right\|_{F}=1.79 \cdot 10^{-11}$ with $\widehat{X} \equiv\left[\begin{array}{lll}\widehat{x}_{1} & \ldots & \widehat{x}_{13}\end{array}\right]$.

$$
n(0,8)=v_{-}\left(S_{11}^{\alpha}\right)-v_{+}\left(Z_{N}^{T} K_{G} Z_{N}\right)=13-0=13
$$

This also matches the number of computed eigenvalues in the interval.

## 6 | CONCLUDING REMARK

We studied the buckling eigenvalue problem of singular pencil, and addressed two open issues associated with the shift-invert Lanczos method. We found that the proposed scheme for counting the number of eigenvalues is a reliable tool for validation.

It is still an open problem how to choose the positive definite matrices $H_{N}$ and $H_{C}$ for the optimal condition number $\kappa_{2}(M)$. An analysis following the work in Reference 25 is a direction of future research.

Also note that there are different implementations of the matrix-vector product $u=C v$. Similar validation schemes can be developed. Performance of different implementations for practical industrial examples is a subject of further study.

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## DATA AVAILABILITY STATEMENT

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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## APPENDIX A. CANONICAL FORM OF A SYMMETRIC SEMI-DEFINITE PENCIL $\boldsymbol{A}-\lambda B$

In this section, we give a constructive derivation of a canonical form of a symmetric semi-definite pencil $A-\lambda B$, namely $A$ is symmetric and $B$ is symmetric semi-positive definite.
Theorem 7. For a symmetric semi-definite pencil $A-\lambda B$, there exists a non-singular matrix $W \in \mathbb{R}^{n \times n}$ such that

$$
W^{T} A W=\begin{align*}
& 2 n_{0}  \tag{A1}\\
& n_{1} \\
& n_{2} \\
& n_{3}
\end{align*}\left[\begin{array}{lllll}
n_{3} n_{0} & n_{1} & n_{2} & n_{3} \\
S & & & \\
& \Lambda_{1} & & \\
& & \Lambda_{2} & \\
& & & 0
\end{array}\right] \text { and } W^{T} B W=\begin{array}{llll}
2 n_{0} \\
n_{1} \\
n_{2} \\
n_{3}
\end{array}\left[\begin{array}{llll}
2 n_{0} & n_{1} & n_{2} & n_{3} \\
\Omega & & & \\
& I_{n_{1}} & \\
& & 0 & \\
& & & 0
\end{array}\right] \text {, }
$$

where

$$
S \equiv I_{n_{0}} \otimes\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right], \quad \Omega \equiv I_{n_{0}} \otimes\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right]
$$

$\Lambda_{1}$ and $\Lambda_{2}$ are diagonal matrices with real diagonal entries, and $\Lambda_{2}$ is nonsingular. Moreover, we have

$$
\begin{aligned}
& n_{0}=\operatorname{dim}(\mathcal{N}(B))-n_{2}-n_{3}, \\
& n_{1}=\operatorname{rank}(B)-n_{0}, \\
& n_{2}=\operatorname{rank}\left(\mathcal{P}_{\mathcal{N}(B)} A \mathcal{P}_{\mathcal{N}(B)}\right), \\
& n_{3}=\operatorname{dim}(\mathcal{N}(A) \cap \mathcal{N}(B)),
\end{aligned}
$$

where $\mathcal{P}_{\mathcal{N}(B)}$ is the orthogonal projection onto $\mathcal{N}(B)$.
We first introduce the following lemma due to Fix and Heiberger, ${ }^{27}$ also see Reference 26, section 15.5.
Lemma 5. For the symmetric semi-definite pencil $A-\lambda B$, there exists a non-singular matrix $W \in \mathbb{R}^{n \times n}$ such that

$$
W^{T} A W=\begin{array}{r}
n_{0} \\
n_{1} \\
n_{1} \\
n_{0} \\
n_{3}
\end{array}\left[\begin{array}{rrrrr}
n_{0} & n_{1} & n_{2} & n_{0} & n_{3} \\
A_{00} & A_{01} & A_{02} & \Sigma & 0 \\
A_{01}^{T} & A_{11} & A_{12} & & \\
A_{02}^{T} & A_{12}^{T} & \Lambda_{2} & & \\
\Sigma & & & 0 & \\
0 & & & & 0
\end{array}\right] \text { and } W^{T} B W \begin{array}{lllll}
n_{0} \\
n_{1} \\
n_{2}
\end{array}\left[\begin{array}{lllll}
n_{0} & n_{n_{0}} & & & n_{2} \\
& n_{0} & n_{3} \\
& I_{n_{1}} & & \\
& & 0 & \\
& & & \\
n_{0} \\
& & & 0 & \\
& & & & 0
\end{array}\right],
$$

where $\Lambda_{2}$ and $\Sigma$ are non-singular, diagonal matrices with real diagonal entries.

Proof. Proof of Theorem 7. By Lemma 5, there exists a non-singular matrix $W_{0} \in \mathbb{R}^{n \times n}$ such that
where $\Lambda_{2}$ and $\Sigma$ are nonsingular, diagonal matrices with real diagonal entries.

Let

$$
W_{1} \equiv \begin{gathered}
\\
n_{0} \\
n_{1} \\
n_{2} \\
n_{0} \\
n_{n_{0}}
\end{gathered}\left[\begin{array}{ccccc}
n_{0} & n_{1} & n_{2} & n_{0} & n_{3} \\
& I_{n_{1}} & & & \\
& & I_{n_{2}} & & \\
-\Sigma^{-1} A_{00} / 2 & -\Sigma^{-1} A_{01} & -\Sigma^{-1} A_{02} & I_{n_{0}} & \\
& & & & I_{n_{3}}
\end{array}\right]
$$

then

Next let

$$
W_{2} \equiv \begin{array}{ccccc} 
\\
n_{0} \\
n_{1} \\
n_{0}
\end{array}\left[\begin{array}{ccccc}
n_{0} & n_{1} & n_{2} & n_{0} & n_{3} \\
I_{n_{0}} & & & & \\
& I_{n_{1}} & & & \\
& -\Lambda_{2}^{-1} A_{12}^{T} & I_{n_{2}} & & \\
& & & I_{n_{0}} & \\
& & & & I_{n_{3}}
\end{array}\right],
$$

then
where $C_{11} \in \mathbb{R}^{n_{1} \times n_{1}}$ is symmetric and $C_{11}=A_{11}-A_{12} \Lambda_{2}^{-1} A_{12}^{T}$.
Define the permutation matrix

$$
P_{3} \equiv\left[\begin{array}{ccccc}
I_{n_{0}} & 0 & 0 & 0 & 0 \\
0 & 0 & I_{n_{1}} & 0 & 0 \\
0 & 0 & 0 & I_{n_{2}} & 0 \\
0 & I_{n_{0}} & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & I_{n_{3}}
\end{array}\right],
$$

then

Since $C_{11} \in \mathbb{R}^{n_{1} \times n_{1}}$ is symmetric, it admits the eigen-decomposition

$$
C_{11}=Q_{1} \Lambda_{1} Q_{1}^{T}
$$

where $Q_{1} \in \mathbb{R}^{n_{1} \times n_{1}}$ is an orthogonal matrix and $\Lambda_{1} \in \mathbb{R}^{n_{1} \times n_{1}}$ is a diagonal matrix. Applying the congruent transformation associated with $W_{4} \equiv \operatorname{diag}\left(I_{n_{0}}, \Sigma^{-1}, Q_{1}, I_{n_{2}}, I_{n_{3}}\right)$, we have

Last, define the permutation matrix $P_{5} \equiv \operatorname{diag}\left(E, I_{n_{1}}, I_{n_{2}}, I_{n_{3}}\right)$ with $E \equiv\left[\begin{array}{llll}e_{1} & e_{n_{0}+1} & e_{2} & \ldots\end{array} e_{2 n_{0}}\right]$ and we have the canonical form in (A1)
where

$$
S \equiv I_{n_{0}} \otimes\left[\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right] \quad \text { and } \quad \Omega \equiv I_{n_{0}} \otimes\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right]
$$

The canonical form (A1) is obtained with $W \equiv W_{0} W_{1} W_{2} P_{3} W_{4} P_{5}$.
Now we interpret the dimension of each block matrix. From the canonical form of $B$ in Equation (A1), we can infer that $n_{0}=\operatorname{dim}(\mathcal{N}(B))-n_{2}-n_{3}$ and $n_{1}=\operatorname{rank}(B)-n_{0}$. Also, $n_{3}=\operatorname{dim}(\mathcal{N}(A) \cap \mathcal{N}(B))$. To interpret $n_{2}$, let $Z \in \mathbb{R}^{n \times\left(n_{0}+n_{2}+n_{3}\right)}$ be the basis of $\mathcal{N}(B)$ consisting of the columns of $W$ and consider the QR decomposition of $Z=Q R$. Since $Q$ is an orthonormal basis of $\mathcal{N}(B), \operatorname{rank}\left(\mathcal{P}_{\mathcal{N}(B)} A \mathcal{P}_{\mathcal{N}(B)}\right)=\operatorname{rank}\left(Q^{T} A Q\right)$. By the Sylvester's law, $\operatorname{rank}\left(Q^{T} A Q\right)=\operatorname{rank}\left(Z^{T} A Z\right)$. But, from the canonical form (A1), $Z^{T} A Z=\operatorname{diag}\left(0_{n_{0}}, \Lambda_{2}, 0_{n_{3}}\right)$ and $\operatorname{rank}\left(Z^{T} A Z\right)=n_{2}$. Therefore, $n_{2}=\operatorname{rank}$ $\left(\mathcal{P}_{\mathcal{N}(B)} A \mathcal{P}_{\mathcal{N}(B)}\right)$.
Corollary 1. The symmetric semi-definite pencil $A-\lambda B$ is simultaneously diagonalizable if and only if $n_{0}=0$. In this case, we have the canonical form

$$
W^{T} A W=\begin{array}{ccc}
n_{1} & n_{2} & n_{3} \\
n_{2} \\
n_{3}
\end{array}\left[\begin{array}{lll}
\Lambda_{1} & & \\
& \Lambda_{2} & \\
& & 0
\end{array}\right] \text { and } \begin{array}{rll}
n_{1} & n_{2} & n_{3} \\
& & W^{T} B W= \\
n_{2} \\
n_{3}
\end{array}\left[\begin{array}{lll}
I_{n_{1}} & & \\
& 0 & \\
& & 0
\end{array}\right]
$$

Proof. From the pairs $(S, \Omega)$ and $\left(\Lambda_{2}, 0\right)$ in Equation (A1), we note that the algebraic and geometric multiplicity of the infinite eigenvalues are $2 n_{0}+n_{2}$ and $n_{0}+n_{2}$, respectively. Therefore, the symmetric semi-definite pencil $A-\lambda B$ is simultaneously diagonalizable if and only if $n_{0}=0$.

